Characterization of the Magnetic Phase in Ti-Doped Vanadium Dioxide M. R. Goeks^{*a}, P. W. Mengyan^{a,b}, R. L. Lichti^b

Project Purpose:

Use μ^+ as experimentally accessible analog to H and probe of the local magnetic environment in Vanadium Dioxide (VO₂) to:

- Identify transition mechanism (*Peierls* vs *Mott-Hubbard*)
- Understand *role* dopants play in changing material properties such as transitions, conductivity, bandgap, magnetism, *etc*
- Explore magnetism caused by disruption of V-V dimer
- Study H behavior (stability/dynamics, energy barriers, *etc*)

VO₂ Properties [1-4]:

- Reversible, Metal-Semiconductor Transition ($T_{MST} = 340K$) triggered thermally, optically, electrically, barometrically
- Metallic (T >T_{MST}, E_G \sim 0eV): Rutile (*Fig 1*), reflective, $\sigma \sim 10^3 - 10^4 (\Omega \cdot \text{cm})^{-1}$
- Semiconducting (T< T_{MST}, E_G \sim 1eV): Monoclinic (*Fig 2*), translucent, $\sigma \sim 10^{-1} - 10^{-3} (\Omega \cdot \text{cm})^{-1}$
- Dopants modify transition temperature and properties
 - Ti, W, Au (etc.): lowers T_{MST} ; F, Cr, Al (etc.): raises T_{MST} [3] minimal effects on properties other than T_{MST}
 - Adding H reduces T_{MST} while increasing σ in monoclinic phase with 3.8 at% H, $T_{MST} \sim 200$ K with nominal σ difference between phases. *Effect* is known but *role* H plays is unknown. [4]





Fig. 2: Monoclinic structure [2]

Fig. 1: Rutile structure [2] where V^{4^+} has loosely bound $e^- \rightarrow high \sigma$

Project Significance:

V pair with and without twisting. *Pairs* e^{-} *into singlet state* $\rightarrow low \sigma$

This work directly contributes to understanding the role and behavior of H in bulk VO₂ compounds; it provides the initial discovery and characterization of the magnetic phases and provides insight into the debated mechanism responsible for the MST.

VO₂ has many potential applications such as in smart windows, smart radiators, microwave wave guides, optical fiber, ultrafast tunable antenna [7]. Applications require exposure to H, which has a major effect on the transitions and electrical properties [4]. The *effect* of H is known but the *role* H plays in modifying the host has yet to be understood.

For applications, it is essential to understand aspects such as how (1) H may propagate into and behave within VO_2

(2) magnetic and electrical properties behave in the compounds

References:

- [1] e.g.: F.J. Morin, *Phys Rev Lett* **3** (1959) 34; A. Cavalleri *et al. J Phys Soc Jpn* **75** (2006) 011004; M. Imada et al. Rev Mod Phys 70 (1998) 1039; J.B. Goodenough, J. Solid State Chem 3 (1971) 490.

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Here we present an initial characterization of a substantial mag B.J. Kim et al. Appl Phys Lett 90 (2007) 023515; M. M. Qazilbash et al. Appl Phys Lett 92 (2008) 241906; phase in three Ti-doped samples of VO₂ by way of Zero Field [2] M. Nazarri, PhD Dissertation, Texas Tech, 2013. Spin Rotation and Relaxation measurements. Our results are [3] i.e.: P. Kiri et al. A dv Mat Lett 1 (2010) 86; Burkhardt et al. Thin Solid Films 345 (1999) 229; commensurate with our related work on a W-doped VO₂ comp [4] C. Wu et al. J Am Chem Soc 133 (2011) 13798. where we see similar shifts in the onset of the magnetic phase [5] A. Yaouanc & P.D. de Reotier, *Muon Spin Rotation* [...] (Oxford Press, New York, 2011); suggesting that the spatial distortion in the structure due to the J.H. Brewer in Encyclopedia of Applied Physics, Muon Spin [...], ed. Trigg, (VCH, New York, 1994) V11 presence of the impurities stabilizes the magnetic phase to high [6] R.S. Hayano *et al. Phys Rev B* **20** (1979) 850. temperatures. The magnetism itself is introduced by a disruptic [7] A. Cavalleri et al. J. Phys Soc Japan 75 (2006) 011004; B. J. Kim et al. Appl Phys Lett 90 (2007) 023515; V-V dimerization and is observed to occur in stoichiometric V around 35K. [**] Funding by: Welch Foundation (grant D-1321), NMU Freshman Fellows program, NMU Foundation [***] Special thanks to the MuSR staff scientists and technicians at TRIUMF and RAL

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Muon Spin Relaxation (MuSR) [5]:



- Implant μ^+ into sample
- μ^+ starts clock
- μ^+ precesses in $B_{eff} = B_{ext} + B_{int}$
- μ^+ decays, emits e⁺ preferentially along direction of spin
- e⁺ hits detector, stops clock
- Time evolution of μ^+ spin tracked via A(t) of e^+ emission

Fig. 3: Experimental setup for a typical MuSR experiment with an externally applied field [5]. Only one of four positron detectors shown.



Fits to predominant frequency in ZF time domain spectra directly measure magnetic field at μ^+ site, yield • 5 at%: $B_{loc} = 111.8 \pm 0.5 \text{ mT as } T \rightarrow 0$; $T_C = 175 \pm 2 \text{ K}$

- 3 at%: $B_{loc} = 224.7 \pm 0.5 \text{ mT as } T \rightarrow 0$; $T_{C} = 80.0 \pm 1 \text{ K}$
- 1 at%: $T_C = -37 \pm 2K$ [not shown above]. B_{loc} is clearly present and develops with decreasing T, however, below 20K significant changes in the time domain signal indicate multiple fields present with very broad distribution.

Conclusion:

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The Experiment:

- ZF (B_{ext}=0T) MuSR Measurements
- *HiTime* and *Helios* spectrometers, M15 & M20C surface muon channels at TRIUMF (Vancouver, Canada)
- Bulk sintered $7 \times 7 \times 1$ mm³ samples, VO₂:Ti (1, 3, 5 at%)
- Temperature range: ~2K to ~285K



Fig. 4: Sample mounted for use in HiTime surrounded by 4 detectors (scintillator cubes). The sample is covered by a thin piece of foil tape to hold it in position directly on top of another scintillator (veto cube) that is used to detect any μ^{-} that miss the sample so that only events from muon decay inside the sample are recorded.

Time domain asymmetry spectra evolve with temperature (Fig 4) and are characterized here by the model:

$$\begin{aligned} (t) &= \sum_{i} A_{i}e^{-\lambda_{i}t}\cos(\gamma_{\mu}B_{eff,i}t + \phi_{i}) \\ &+ \sum_{i} A_{j}\left(\frac{1}{3} + \frac{2}{3}\left(1 - \Delta_{j}^{2}t^{2}\right)e^{-\frac{\Delta_{j}^{2}t^{2}}{2}}\right) \\ &+ \sum_{i} A_{j}\left(\frac{1}{3} + \frac{2}{3}\left(1 - \Delta_{j}^{2}t^{2}\right)e^{-\frac{\Delta_{j}^{2}t^{2}}{2}}\right) \\ &- \text{Static Kubo-Toyabe (sKT) component} \\ &- \Delta_{j} = \text{Second moment of measured field dist} \end{aligned}$$

Fig 4 [left]: Example of time domain spectra from VO_2 : Ti 3 at%, at 2.4 K, 72.7 K, and 77.6 K. (a) The early time shows a fast relaxing oscillation, indicating a well-defined magnetic phase which is present only below Tc. (b) Full time spectra showing major differences in initial asymmetry above and below T_C , typical for a magnetic transition. The highest temperature spectra shown is typical for a state where only nuclear moments contribute to the field sensed by the μ^+ and is characterized by a Kubo-Toyabe function [6].



Fig 7: Temperature dependence of compo*nent asymmetry for VO*₂:*Ti 3 at%. Asymmetry* is proportional to the fraction of the material in that state. As T approaches T_C , the oscillat*ing component (magnetic phase,* **■***) is re*placed by a non-relaxing component () likely from the increasingly fast fluctuations associated with the break-up of the low T long range order. At T_C , the oscillation is complete*ly replaced by the static Kubo-Toyabe* (sKT,•) indicating only nuclear fields are present.

| | Project Status, Future Work: |
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| gnetic Muon | • A refined model is needed to characterize other fields present (not presented here) in the spectra. These additional fields suggest AFM order or other magnetic phases. |
| oound, | • Analyze Mu (H-like) behavior in the non-magnetic regime to provide insight into how H behaves and what role H impurities play in modifying the properties of VO ₂ . |
| her on of the O_2 | • Combine these results with data from stoichiometric VO ₂ and VO ₂ :W (1 and 2.4 weight %) to establish trends with impurity content to begin to develop a robust model for the transitions (magnetic, structural, optical, electric) in these compounds. |





 VO_2 and impurity insitions